ON SOME NONLINEAR FRACTIONAL EQUATIONS INVOLVING THE BESSEL OPERATOR

SIMONE SECCHI

Dedicated to Francesca

ABSTRACT. Under different assumptions on the potential functions b and c, we study the fractional equation $(I-\Delta)^{\alpha}u=\lambda b(x)|u|^{p-2}u+c(x)|u|^{q-2}u$ in \mathbb{R}^{N} . Our existence results are based on compact embedding properties for weighted spaces.

1. Introduction

In this paper we provide some existence results for a class of nonlinear fractional equations of the form

(1.1)
$$(I - \Delta)^{\alpha} u = \lambda b(x) |u|^{p-2} u + c(x) |u|^{q-2} u in \mathbb{R}^{N}$$

where $0 < \alpha < 1$, p and q belong to the interval $(0, 2^*_{\alpha})$ with $2^*_{\alpha} = 2N/(N-2\alpha)$. We will assume throughout that

(H): The potential functions b and c are continuous and bounded.

Equations involving the nonlocal operator $(I-\Delta)^{\alpha}$ arise in the study of standing waves $\psi=\psi(t,x)$ for Schrödinger–Klein–Gordon equations of the form

$$i\frac{\partial \psi}{\partial t} = (I - \Delta)^{\alpha} \psi - f(x, \psi), \quad (t, x) \in \mathbb{R} \times \mathbb{R}^{N}$$

which describe the behavior of bosons. We refer to [22, 23] for a physical introduction to these fractional equations.

The case $\alpha = 1$ corresponds to the classical Schrödinger equation $-\Delta u + u = f(x, u)$, and we cannot review the huge literature here. In the last years the so-called fractional Laplacian $(-\Delta)^{\alpha}$ has become popular in the community of Nonlinear Analysts: we refer the interested reader to the recent survey [6] and to the references therein.

The most physically relevant fractional case is $\alpha = 1/2$, and the corresponding Bessel operator, or more precisely the operator $\sqrt{-\Delta + m^2} - m$, goes under the name of relativistic Schrödinger operator.

It is interesting to quote a sentence from [7, page 119]:

One of the fundamental mathematical problems in proving the stability (or instability) of matter is to estimate the infimum of the spectrum of the operator $H = H_0 + V$ when the numbers N and/or M become large. In this asymptotic regime, the free Hamiltonian $H_0 = F(p) = \sqrt{-\Delta + m^2} - m$ can be abandoned

Date: June 15, 2015.

 $^{2010\} Mathematics\ Subject\ Classification.\ 35\text{J}60, 35\text{Q}55, 35\text{S}05.$

Key words and phrases. Fractional Sobolev Spaces, Bessel spaces, Fractional Laplacian.

Supported by FIRB 2012 "Dispersive equations: Fourier analysis and variational methods" and by PRIN 2012 "Aspetti variazionali e perturbativi nei problemi differenziali nonlineari".

and replaced by $H_0 = |p|$. Indeed, the difference between these two operators remains bounded and the asymptotic result needed to prove the stability of matter can be proved using either one of these free Hamiltonians. Obviously, the scaling properties of the function F(p) = |p| attracted early investigations of the corresponding pseudo-differential operator H_0 especially because its scaling is related to the scaling of the Coulomb potential. Also, very fine estimates on its Green's function are available. These are the technical reasons why F(p) = |p| is preferred to $\sqrt{p^2 + m^2} - m$.

We will come back to the issue of scaling properties in the last Section. R. Frank *et al.* in [18, page 5] promise to extend part of their spectral theory for the fractional Laplacian to the pseudodifferential operator $(-\Delta + m^2)^{\alpha}$.

The equation $\sqrt{I-\Delta}\ u=f(u)$ was studied in [37] by means of a Dirichlet-to-Neumann local realization that was extended to any $\alpha\in(0,1)$ by Fall *et al.* in [15]. The case $\alpha=1/2$ was also studied in [11–14] with a non-local convolutions term on the right-hand side.

On the contrary, very few papers deal with equation (1.1) for arbitrary $\alpha \in (0,1)$. The very recent paper [17] deals with the case in which the nonlinearity f = f(x,u) is essentially of the form $\bar{f}(u) + a(x)(|u| + |u|^p)$ with $\lim_{|x| \to +\infty} a(x) = 0$, in the spirit of [30].

In this paper we limit ourselves to a somehow particular class of nonlinearities as in (1.1), and we provide a few existence results under different assumptions on the two potential functions b and c

Section 2 contains some preliminaries on the functional setting that we use to solve (1.1). In Section 3 we begin with the case p > 2, q > 2, b = 0 identically and c positive and "vanishing at infinity" in an appropriate sense. In Section 4 we allow b to tend to a positive constant at infinity, and in Section 5 we consider the concave-convex case 1 with sign-changing potentials <math>b and c.

Notation.

- (1) The letter C will stand for a generic positive constant that may vary from line to line.
- (2) The symbol $\|\cdot\|_p$ will be reserved for the norm in $L^p(\mathbb{R}^N)$.
- (3) The operator D will be reserved for the (Fréchet) derivative, also for functions of a single real variable.
- (4) The symbol \mathcal{L}^N will be reserved for the Lebesgue N-dimensional measure.
- (5) The symbol CA will denote the complement of the subset A (usually in \mathbb{R}^N).
- (6) $\mathbb{R}^{N+1}_+ = \{(x,y) \in \mathbb{R}^N \times [0,+\infty) \}$
- (7) For a real-valued function f, we set $f^+ = \max\{f, 0\}$, the positive part of f. The negative part of f is defined similarly.
- (8) The Fourier transform of a function f will be denoted by $\mathcal{F}u$.

2. Preliminaries and functional setting

For $\alpha > 0$ we introduce the Bessel function space

$$L^{\alpha,2}(\mathbb{R}^N) = \left\{ f \colon f = G_\alpha * g \text{ for some } g \in L^2(\mathbb{R}^N) \right\},$$

where the Bessel convolution kernel is defined by

(2.1)
$$G_{\alpha}(x) = \frac{1}{(4\pi)^{\alpha/2} \Gamma(\alpha/2)} \int_0^\infty \exp\left(-\frac{\pi}{t}|x|^2\right) \exp\left(-\frac{t}{4\pi}\right) t^{\frac{\alpha-N}{2}-1} dt$$

The norm of this Bessel space is $||f|| = ||g||_2$ if $f = G_\alpha * g$. The operator $(I - \Delta)^{-\alpha}u = G_{2\alpha} * u$ is usually called Bessel operator of order α .

In Fourier variables the same operator reads

(2.2)
$$G_{\alpha} = \mathcal{F}^{-1} \circ \left(\left(1 + |\xi|^2 \right)^{-\alpha/2} \circ \mathcal{F} \right),$$

so that

$$||f|| = \left| (I - \Delta)^{\alpha/2} f \right|_2.$$

For more detailed information, see [1, 33] and the references therein.

Remark 2.1. In the paper [15] the pointwise formula

(2.3)
$$(I - \Delta)^{\alpha} u(x) = c_{N,\alpha} \, \text{P. V.} \int_{\mathbb{R}^N} \frac{u(x) - u(y)}{|x - y|^{\frac{N+2\alpha}{2}}} K_{\frac{N+2\alpha}{2}}(|x - y|) \, dy + u(x)$$

was derived for functions $u \in C_c^2(\mathbb{R}^N)$. Here $c_{N,\alpha}$ is a positive constant depending only on N and α , P.V. denotes the principal value of the singular integral, and K_{ν} is the modified Bessel function of the second kind with order ν (see [15, Remark 7.3] for more details). Since a closed formula for K_{ν} is unknown, equation (2.3) is not particularly useful for our purposes.

We recall the embedding properties of Bessel spaces (see [16, 33, 35]).

(1) $L^{\alpha,2}(\mathbb{R}^N) = W^{\alpha,2}(\mathbb{R}^N) = H^{\alpha}(\mathbb{R}^N).$ Theorem 2.2.

- (2) If $\alpha \geq 0$ and $2 \leq q \leq 2_{\alpha}^* = 2N/(N-2\alpha)$, then $L^{\alpha,2}(\mathbb{R}^N)$ is continuously embedded into $L^q(\mathbb{R}^N)$; if $2 \leq q < 2_{\alpha}^*$ then the embedding is locally compact. (3) Assume that $0 \leq \alpha \leq 2$ and $\alpha > N/2$. If $\alpha N/2 > 1$ and $0 < \mu \leq \alpha N/2 1$, then $L^{\alpha,2}(\mathbb{R}^N)$ is continuously embedded into $C^{1,\mu}(\mathbb{R}^N)$. If $\alpha N/2 < 1$ and $0 < \mu \leq \alpha N/2$, then $L^{\alpha,2}(\mathbb{R}^N)$ is continuously embedded into $C^{0,\mu}(\mathbb{R}^N)$.

Remark 2.3. Although the Bessel space $L^{\alpha,2}(\mathbb{R}^N)$ is topologically undistinguishable from the Sobolev fractional space $H^{\alpha}(\mathbb{R}^N)$, we will not confuse them, since our equation involves the Bessel norm.

We collect here a couple of technical lemmas taken from the remarkable paper [29].

Lemma 2.4. Let $0 < \alpha < N/2$ and let $u \in L^{\alpha,2}(\mathbb{R}^N)$. Let $\varphi \in C_0^{\infty}(\mathbb{R}^N)$ and for each R > 0 let $\varphi_R(x) = \varphi(R^{-1}x)$. Then $\lim_{R\to 0} \varphi_R u = 0$ in $L^{\alpha,2}(\mathbb{R}^N)$. If, in addition, φ equals one in a neighborhood of the origin, then $\lim_{R\to +\infty} \varphi_R u = 0$ in $L^{\alpha,2}(\mathbb{R}^N)$.

Proof. Since $L^{\alpha,2}(\mathbb{R}^N)$ is equivalent to $H^s(\mathbb{R}^N)$, the proof of [29, Lemma 5] carries over with only minor modifications.

Lemma 2.5. Let $0 < \alpha < N/2$, let Ω be a bounded open subset of \mathbb{R}^N and let $\varphi \in C_0^{\infty}(\mathbb{R}^N)$. Then the commutator $[\varphi, (I - \Delta)^{\alpha/2}]: L^{\alpha,2}(\mathbb{R}^N) \to L^2(\mathbb{R}^N)$ is a compact operator.

Proof. It suffices to remark that the proof of [29, Lemma 6] actually contains a proof of our statement.

Definition 2.6. We say that $u \in L^{\alpha,2}(\mathbb{R}^N)$ is a weak solution to (1.1) if

$$\int_{\mathbb{R}^N} (I - \Delta)^{\alpha/2} u \ (I - \Delta)^{\alpha/2} v \, dx = \int_{\mathbb{R}^N} b(x) |u|^{p-2} uv \, dx + \int_{\mathbb{R}^N} c(x) |u|^{q-2} uv \, dx$$

for all $v \in L^{\alpha,2}(\mathbb{R}^N)$, or, equivalently,

$$\int_{\mathbb{R}^N} (1 + |\xi|^2)^{\alpha} \mathcal{F} u(\xi) \mathcal{F} v(\xi) \, d\xi = \int_{\mathbb{R}^N} b(x) |u|^{p-2} uv \, dx + \int_{\mathbb{R}^N} c(x) |u|^{q-2} uv \, dx$$

It is readily seen that weak solutions to (1.1) correspond to critical points of the differentiable functional $J: L^{\alpha,2}(\mathbb{R}^N) \to \mathbb{R}$ defined by

$$J(u) = \frac{1}{2} \| (I - \Delta)^{\alpha/2} u \|_2^2 - \frac{\lambda}{p} \int_{\mathbb{R}^N} b(x) |u|^p \, dx - \frac{1}{q} \int_{\mathbb{R}^N} c(x) |u|^q \, dx$$

3. Potentials vanishing at infinity

In this section with treat the particular case of equation (1.1) with $b \equiv 0$, or $\lambda = 0$. We can actually deal with a more general model,

$$(3.1) (I - \Delta)^{\alpha} u = c(x) f(u) in \mathbb{R}^{N}$$

Definition 3.1. We say that a function $c \in L^{\infty}(\mathbb{R}^N)$ satisfies the compactness condition (K) if

$$\lim_{r \to +\infty} \int_{A_n \cap \complement B(0,r)} c(x) \, dx = 0$$

uniformly with respect to n whenever $\{A_n\}_n$ is a sequence of Borel sets such that

$$\sup_{n} \mathcal{L}^{N}(A_{n}) < \infty.$$

Following [2], we assume

- (c1): c > 0 everywhere.
- (c2): c satisfies the condition (K).

Our assumptions on the nonlinearity f read as follows: f is continuous and

- (f1): $\limsup_{s\to 0} f(s)/s = 0$.
- (f2): $\limsup_{s \to +\infty} f(s)/s^{2_{\alpha}^*-1} = 0.$
- (f3): The map $s \mapsto s^{-1}f(s)$ is increasing and the primitive F of f satisfies $\lim_{s \to +\infty} F(s)/s^2 = +\infty$.

Remark 3.2. Since $G_{\alpha} > 0$, it follows easily that $(I - \Delta)^{\alpha}$ satisfies a maximum principle. Therefore, we will look for positive solutions and assume that f(s) = 0 for every $s \leq 0$.

The Euler functional associated to (3.1) is $J: L^{\alpha,2}(\mathbb{R}^N) \to \mathbb{R}$, defined by

$$J(u) = \frac{1}{2} \| (I - \Delta)^{\alpha/2} \|_2^2 - \int_{\mathbb{R}^N} c(x) F(u) \, dx.$$

The following lemma can be proved by standard techniques.

Lemma 3.3. The functional J is of class C^1 has the mountain-pass geometry.

As a consequence, there exists (see [8]) a Cerami sequence $\{u_n\}_n \subset L^{\alpha,2}(\mathbb{R}^N)$ for J, i.e. a sequence such that

(3.2)
$$\lim_{n \to +\infty} J(u_n) = c \text{ and } \lim_{n \to +\infty} (1 + ||u_n||) DJ(u_n) = 0,$$

where

$$c_{\rm mp} = \inf_{\gamma \in \Gamma} \sup_{0 < t < 1} J(\gamma(t))$$

and

$$\Gamma = \left\{\gamma \in C([0,1], L^{\alpha,2}\mathbb{R}^N) \mid \gamma(0) = 0, \ J(\gamma(1)) < 0\right\}.$$

In view of Remark 3.2, we may assume that $u_n \geq 0$ for every $n \in \mathbb{N}$.

Definition 3.4. For a positive and measurable function $w: \mathbb{R}^N \to [0, +\infty)$ we denote

$$L^{q}(w d\mathcal{L}^{N}) = \left\{ u \colon \mathbb{R}^{N} \to \mathbb{R} \mid \int |u(x)|^{q} w(x) dx < \infty \right\}$$

Proposition 3.5. The space $L^{\alpha,2}(\mathbb{R}^N)$ is compactly embedded into $L^q(c d\mathcal{L}^N)$ for every $2 < q < 2_{\alpha}^*$.

Proof. We fix $q \in (2, 2^*_{\alpha})$ and pick $\varepsilon > 0$. There exist $s_0 < s_1$ such that

$$c(x)|s|^q \le \varepsilon C\left(s^2 + |s|^{2^*_{\alpha}}\right) + Cc(x)\chi_{[s_0,s_1]}(|s|)|s|^{2^*_{\alpha}}$$

for every $s \in \mathbb{R}$. Integrating,

$$\int_{\complement B(0,r)} c(x) |u(x)|^q \, dx \leq \varepsilon C Q(u) + C \int_{A \cap \complement B(0,r)} c(x) \, dx$$

for every $u \in L^{\alpha,2}(\mathbb{R}^N)$, where $Q(u) = ||u||_2^2 + ||u||_{2^*_{\alpha}}^{2^*_{\alpha}}$ and $A = \{x \in \mathbb{R}^N \mid s_0 \le |u(x)| \le s_1\}$. If $v_n \to 0$ weakly in $L^{\alpha,2}(\mathbb{R}^N)$, then

$$\left\| (I - \Delta)^{\alpha/2} v_n \right\|_2^2 \le M_1 \quad \text{and} \quad \|v_n\|_{2^*_{\alpha}}^{2^*_{\alpha}} \le M_1$$

for some $M_1 > 0$. In particular, the sequence $\{Q(v_n)\}_n$ is bounded in \mathbb{R} . On the other hand, if we set

$$A_n = \left\{ x \in \mathbb{R}^N \mid s_0 \le |v_n(x)| \le s_1 \right\},\,$$

then

$$s_0^{2_{\alpha}^*} \mathcal{L}^N(A_n) \le \int_{A_n} |v_n|^{2_{\alpha}^*} dx \le M_1$$

for all $n \in \mathbb{N}$. This shows that $\sup_n \mathcal{L}^N(A_n) < \infty$ and from our assumption on c there exists a number r > 0 such that

$$\int_{A_n\cap \mathbb{C}B(0,r)} c(x)\,dx < \frac{\varepsilon}{s_1^{2_\alpha^*}}.$$

We are ready to conclude that

$$\int_{B(0,r)} c(x) |v_n(x)|^q \, dx \le \varepsilon C M_1 + s_1^{2_\alpha^*} \int_{A_n \cap \mathcal{L}B(0,r)} c(x) \, dx < (C M_1 + 1)\varepsilon$$

for all $n \in \mathbb{N}$. From the compact embedding of $L^{\alpha,2}(\mathbb{R}^N)$ into $L^q_{loc}(\mathbb{R}^N)$ and the boundedness of c, we know that

$$\lim_{n \to +\infty} \int_{B(0,r)} c(x) |v_n(x)|^q dx = 0.$$

Hence $v_n \to 0$ strongly in $L^q(c d\mathcal{L}^N)$. The proof is complete.

Remark 3.6. If assumption (c2) is replaced by $\lim_{|x|\to+\infty} c(x) = 0$, then the compactness of the embedding is classical: see for instance [36, Lemma 3.2]. A different assumption ensuring the compactness of the weighted embedding appears in [9].

Proposition 3.7. If $v_n \to v$ weakly in $L^{\alpha,2}(\mathbb{R}^N)$ then

(3.3)
$$\lim_{n \to +\infty} \int_{\mathbb{R}^N} c(x) F(v_n(x)) dx = \int_{\mathbb{R}^N} c(x) F(v(x)) dx$$

and

(3.4)
$$\lim_{n \to +\infty} \int_{\mathbb{R}^N} c(x) f(v_n(x)) v_n(x) dx = \int_{\mathbb{R}^N} c(x) f(v(x)) v(x) dx$$

Proof. Let us fix $q \in (2, 2^*_{\alpha})$ and $\varepsilon > 0$. From (f1)–(f3) there exists a constant C > 0 such that

$$|c(x)F(s)| \le \varepsilon C\left(s^2 + |s|^{2^*_{\alpha}}\right) + c(x)|s|^q$$

for every $s \in \mathbb{R}$ and every $x \in \mathbb{R}^N$. Proposition 3.5 tells us that

$$\lim_{n \to +\infty} \int_{\mathbb{R}^N} c(x) |v_n(x)|^q dx = \int_{\mathbb{R}^N} c(x) |v(x)|^q dx,$$

so that for some r > 0,

$$\int_{\mathsf{C}B(0,r)} c(x) |v_n(x)|^q \, dx < \varepsilon \quad \text{for all } n \in \mathbb{N}.$$

But $\{v_n\}_n$ is bounded in $L^{\alpha,2}(\mathbb{R}^N)$ and therefore also in $L^2(\mathbb{R}^N)$ and in $L^{2^*_{\alpha}}(\mathbb{R}^N)$. We choose $M_2 > 0$ such that $\int_{\mathbb{R}^N} v_n(x)^2 dx \leq M_2$ and $\int_{\mathbb{R}^N} |v_n(x)|^{2^*_{\alpha}} dx \leq M_2$ for all $n \in \mathbb{N}$. Integrating (3.5) we get

$$\left| \int_{\mathbb{C}B(0,r)} c(x) F(v_n(x)) \, dx \right| < (2CM_2 + 1) \varepsilon \quad \text{for all } n \in \mathbb{N}.$$

Applying a general convergence result (see for instance [10, Lemma 2.4]) to the sequence $\{\|v_n\|_{L^{2_{\alpha}^*}(B(0,r))}\}_n$ together with our assumption (f2) we conclude that

$$\lim_{n \to +\infty} \int_{B(0,r)} c(x) F(v_n(x)) \, dx = \int_{B(0,r)} c(x) F(v(x)) \, dx.$$

This shows the validity of (3.3), and a similar argument proves also (3.4).

Proposition 3.8. The sequence $\{u_n\}_n$ introduced in (3.2) is bounded in $L^{\alpha,2}(\mathbb{R}^N)$.

Proof. For every $n \in \mathbb{N}$, let $t_n \in [0,1]$ be chosen so that $J(t_n u_n) = \max_{0 \le t \le 1} J(t u_n)$. Let us prove that the sequence $\{J(t_n u_n)\}_n$ is bounded from above in \mathbb{R} . The conclusion is trivial if either $t_n = 0$ or $t_n = 1$. If $0 < t_n < 1$, then $DJ(t_n u_n)u_n = 0$. As a consequence,

$$2J(t_n u_n) = 2J(t_n u_n) - DJ(t_n u_n) u_n = \int_{\mathbb{R}^N} c(x) H(t_n u_n(x)) dx,$$

for H(s) = sf(s) - 2F(s). Since $u_n \ge 0$ and H is non-decreasing,

$$2J(t_n u_n) \le \int_{\mathbb{R}^N} c(x)H(u_n(x)) dx = 2J(u_n) - DJ(u_n)u_n = 2J(u_n) + o(1).$$

This shows that $\{J(t_nu_n)\}_n$ is bounded from above. To complete the proof, we argue by contradiction. Let us assume that (possibily along a subsequence) $\lim_{n\to+\infty} \|(I-\Delta)^{\alpha/2}u_n\|_2 = +\infty$. We normalize u_n by introducing $w_n = u_n/\|(I-\Delta)^{\alpha/2}u_n\|_2$. Without loss of generality, we may assume that $w_n \to w$ weakly in $L^{\alpha,2}(\mathbb{R}^N)$. Divinding out the relation $J(u_n) \to c$, we find

$$o(1) + \frac{1}{2} = \int_{\mathbb{R}^N} \frac{c(x)F(u_n(x))}{\|(I - \Delta)^{\alpha/2}u_n\|_2^2} dx = \int_{\mathbb{R}^N} \frac{c(x)F(u_n(x))}{|u_n(x)|^2} w_n(x)^2 dx.$$

By (f3), to each T > 0 we can attach $\xi > 0$ with the property that $|s| \ge \xi$ implies $F(s) \ge Ts^2$. Therefore

$$o(1) + \frac{1}{2} \ge \int_{\mathbb{C}w^{-1}(0) \cap \{|u_n| > \xi\}} \frac{c(x)F(u_n(x))}{|u_n(x)|^2} w_n(x)^2 dx \ge T \int_{\mathbb{C}w^{-1}(0) \cap \{|u_n| > \xi\}} c(x)w_n(x)^2 dx.$$

An application of Fatou's lemma yields $\frac{1}{2} \geq T \int_{\mathbb{C}w^{-1}(0)} c(x)w(x)^2 dx$ and consequently w = 0 by letting $T \to +\infty$.

Finally, given T > 0, we have

$$J(t_n u_n) \ge J\left(\frac{T}{\|(I - \Delta)^{\alpha/2} u_n\|_2} u_n\right) = J(Tw_n) = \frac{T^2}{2} - \int_{\mathbb{R}^N} c(x) F(Tw_n(x)) dx.$$

Proposition 3.7 yields $\lim_{n\to+\infty}\int_{\mathbb{R}^N}c(x)F(Tw_n(x))\,dx=0$ and thus

$$\liminf_{n \to +\infty} J(t_n u_n) \ge \frac{T^2}{2}.$$

Letting $T \to +\infty$ we contradict the boundedness from above of the sequence $\{J(t_n u_n)\}_n$. The proof is complete.

Let us state and prove the main result of this section.

Theorem 3.9. Suppose that (f1)–(f3) and (c1)–(c2) hold. Then equation (3.1) possesses at least one positive solution. Moreover $u \in C^{0,\mu}(\mathbb{R}^N)$ for some $q_0 \in [2,+\infty)$ and $\mu \in (0,1)$.

Proof. We first prove that our Cerami sequence $\{u_n\}_n$ is relatively compact in $L^{\alpha,2}(\mathbb{R}^N)$. To this aim, by virtue of Proposition 3.8, we may assume that, up to a subsequence, $u_n \to u$ weakly in $L^{\alpha,2}(\mathbb{R}^N)$. Since $\{u_n\}_n$ is a Cerami sequence, we find

$$\lim_{n \to +\infty} \left\| (I - \Delta)^{\alpha/2} \right\|_2^2 = \lim_{n \to +\infty} \int_{\mathbb{R}^N} c(x) f(u_n(x)) u_n(x) \, dx.$$

By Proposition 3.7,

$$\lim_{n \to +\infty} \int_{\mathbb{R}^N} c(x) f(u_n(x)) u_n(x) dx = \int_{\mathbb{R}^N} c(x) f(u(x)) u(x) dx.$$

Exploiting again the fact that $DJ(u_n)u \to 0$, we have

$$\left\| (I - \Delta)^{\alpha/2} u \right\|_2^2 = \int_{\mathbb{R}^N} c(x) f(u(x)) u(x) \, dx.$$

In particular $\lim_{n\to+\infty} \|(I-\Delta)^{\alpha/2}u_n\|_2^2 = \|(I-\Delta)^{\alpha/2}u\|_2^2$, or $u_n\to u$ strongly in $L^{\alpha,2}(\mathbb{R}^N)$. But $J(u)=c_{\mathrm{mp}}$ and DJ(u)=0, so that u weakly solves (3.1). The positivity of u follows from the fact that $u_n\geq 0$ and the positivity of the Bessel function G_α see [17, Proposition 3.2]. The regularity of u follows with minor changes from the arguments developed in [16].

Remark 3.10. We have been sketchy about the regularity theory for our solutions, since the Bessel operator is precisely the main tool to develop a regularity theory for the fractional Laplacian, see [16, Appendix A]. In this sense, the fractional Laplacian is *harder* to analyze. Moreover, by similar arguments as those in [17] it can be shown that our solution decays exponentially fast at infinity. This is a common feature for local elliptic partial differential operators, while it is false for the fractional Laplacian, see [16].

4. Potentials having a finite limit

The main tool that we used to solve equation (3.1) is the compactness of the embedding $L^{\alpha,2}(\mathbb{R}^N) \to L^q(c\,d\mathcal{L}^N)$ stated in Proposition 3.5. In this section we study a model case in which a different approach must be used. We consider

$$(4.1) (I - \Delta)^{\alpha} u = \lambda b(x) |u|^{p-2} u + c(x) |u|^{q-2} u \text{in } \mathbb{R}^N,$$

where $p, q \in (2, 2_{\alpha}^*)$ and $\lambda > 0$ is a parameter. We will assume that c > 0 satisfies the condition (K), and impose the following ones on the potential b:

(b1)
$$b \in L^{\infty}(\mathbb{R}^N)$$
, $b \ge 0$ but not identically zero, and $\lim_{|x| \to +\infty} b(x) = \bar{b}$.

Weak solutions to equation (4.1) correspond to critical points of the Euler functional $J_{\lambda} : L^{\alpha,2}(\mathbb{R}^N) \to \mathbb{R}$ defined by

$$J_{\lambda}(u) = \frac{1}{2} \left\| (I - \Delta)^{\alpha/2} u \right\|_{2}^{2} - \frac{\lambda}{p} \int_{\mathbb{R}^{N}} b(x) |u(x)|^{p} dx - \frac{1}{q} \int_{\mathbb{R}^{N}} c(x) |u(x)|^{q} dx.$$

Remark 4.1. The functional $u \mapsto \int c(x)|u(x)|^q dx$ is weakly sequentially continuous by the results of the previous section.

Let us introduce the artificial constraint

$$\mathcal{M}_{\lambda} = \left\{ u \in L^{\alpha,2}(\mathbb{R}^N) \setminus \{0\} \mid DJ_{\lambda}(u)u = 0 \right\},\,$$

and standard arguments show that \mathcal{M}_{λ} is a natural constraint for J_{λ} . In particular, any solution of the minimization problem

$$I_{\lambda} = \inf_{u \in \mathcal{M}_{\lambda}} J_{\lambda}(u)$$

is a solution to equation (4.1).

For $\lambda = 0$, we introduce the constraint

$$\mathcal{M}_0 = \left\{ u \in L^{\alpha,2}(\mathbb{R}^N) \setminus \{0\} \mid DJ_0(u)u = 0 \right\}$$

corresponding to the Euler functional

$$J_0(u) = \frac{1}{2} \left\| (I - \Delta)^{\alpha/2} u \right\|_2^2 - \frac{1}{q} \int_{\mathbb{R}^N} c(x) |u(x)|^q dx.$$

By Remark 4.1, the minimization problem

$$I_0 = \inf_{u \in \mathcal{M}_0} J_0(u)$$

is solved by some function $u_0 \in L^{\alpha,2}(\mathbb{R}^N)$ that satisfies $(I - \Delta)^{\alpha}u_0 = c(x)|u_0|^{q-2}u_0$.

Lemma 4.2. There results $I_{\lambda} \leq I_0$ for all $\lambda > 0$.

Proof. Let $u \in L^{\alpha,2}(\mathbb{R}^N)$ be such that

$$\left\| (I - \Delta)^{\alpha/2} u \right\|_2^2 = \int_{\mathbb{R}^N} c(x) |u(x)|^q dx.$$

Pick $\bar{\sigma} \in (0,1)$ such that $v = \bar{\sigma}u \in \mathcal{M}_{\lambda}$. If we differentiate

$$h(\sigma) = \frac{\sigma^2}{2} \left\| (I - \Delta)^{\alpha/2} u \right\|_2^2 - \frac{\sigma^q}{q} \int_{\mathbb{R}^N} c(x) |u(x)|^q dx$$

and remark that $Dh(\sigma) > 0$ for every $\sigma \in (0,1)$, we may conclude that

$$\begin{split} I_{\lambda}(v) &= \frac{\bar{\sigma}}{2} \left\| (I - \Delta)^{\alpha/2} u \right\|_{2}^{2} - \frac{\bar{\sigma}^{p}}{p} \int_{\mathbb{R}^{N}} \lambda b(x) |u(x)|^{p} dx - \frac{\bar{\sigma}^{q}}{q} \int_{\mathbb{R}^{N}} c(x) |u(x)|^{q} dx \\ &< \frac{\bar{\sigma}}{2} \left\| (I - \Delta)^{\alpha/2} u \right\|_{2}^{2} - \frac{\bar{\sigma}^{q}}{q} \int_{\mathbb{R}^{N}} c(x) |u|^{q} dx \\ &< \frac{1}{2} \left\| (I - \Delta)^{\alpha/2} u \right\|_{2}^{2} - \frac{1}{q} \int_{\mathbb{R}^{N}} c(x) |u|^{q} dx = I_{0}. \end{split}$$

Let us introduce

$$\alpha_1 = \inf_{\|u\|_p = 1} \|(I - \Delta)^{\alpha/2} u\|_2^2 > 0.$$

The main result of this section reads as follows.

Theorem 4.3. Under our assumptions on b and c, if

$$(4.2) I_{\lambda} < \frac{p-2}{2p} \alpha_1^{\frac{p}{p-2}} \left(\lambda \bar{b}\right)^{\frac{2}{2-p}},$$

then equation (4.1) possesses a nontrivial solution.

Proof. We follow the ideas developed in [9]. By Ekeland's variational principle, there exists a minimizing sequence $\{u_n\}_n \subset \mathcal{M}_{\lambda}$ such that $\lim_{n \to +\infty} J_{\lambda}(u_n) = I_{\lambda}$ and $\lim_{n \to +\infty} DJ_{\lambda}(u_n) = 0$. It is readily seen that the sequence $\{u_n\}_n$ is bounded in $L^{\alpha,2}(\mathbb{R}^N)$, so that we may assume without loss of generality that $u_n \to u$ weakly in $L^{\alpha,2}(\mathbb{R}^N)$. Let us set

$$\alpha_{\infty} = \lim_{R \to +\infty} \limsup_{n \to +\infty} \int_{\mathbb{C}B(0,R)} |u_n(x)|^p dx$$
$$\beta_{\infty} = \lim_{R \to +\infty} \limsup_{n \to +\infty} \int_{\mathbb{C}B(0,R)} |(I - \Delta)^{\alpha/2} u_n|^2 dx.$$

We claim that

$$(4.3) \alpha_1 \alpha_{\infty}^{2/p} \le \beta_{\infty}$$

(4.4)
$$\limsup_{n \to +\infty} \int_{\mathbb{R}^N} |u_n(x)|^p dx = \int_{\mathbb{R}^N} |u(x)|^p dx + \alpha_{\infty}$$

(4.5)
$$\limsup_{n \to +\infty} \int_{\mathbb{R}^N} |(I - \Delta)^{\alpha/2} u_n|^2 dx \ge \int_{\mathbb{R}^N} |(I - \Delta)^{\alpha/2} u|^2 dx + \beta_{\infty}.$$

Let us fix a smooth, positive cutoff function φ that equals one on a neighborhood of the origin. Define $\varphi_R = \varphi(\cdot/R)$ and $\tilde{\varphi}_R = 1 - \varphi_R$. By definition of α_1 , we have

$$\alpha_1 \|\tilde{\varphi}_R u_n\|_p^2 \le \|(I - \Delta)(\tilde{\varphi}_R u_n)\|_2^2$$

By Lemmas 2.4 and 2.5 in the Appendix,

$$(I - \Delta)^{\alpha/2} (\tilde{\varphi}_R u_n) = (I - \Delta)^{\alpha/2} u_n - (I - \Delta)^{\alpha/2} (\varphi_R u_n) = (I - \Delta)^{\alpha/2} u_n - \varphi_R (I - \Delta)^{\alpha/2} u_n + o_n(1) = (1 - \varphi_R) (I - \Delta)^{\alpha/2} u_n + o_n(1)$$

where $o_n(1) \to 0$ as $n \to +\infty$ in $L^2(\mathbb{R}^N)$. Therefore

$$\|\alpha_1\|\tilde{\varphi}_R u_n\|_p^p \le \int_{\mathbb{R}^N} |\tilde{\varphi}_R|^2 |(I-\Delta)^{\alpha/2} u_n|^2 dx + o_n(1)$$

Letting $n \to +\infty$ and $R \to +\infty$, we get (4.3). The relations (4.4) and (4.5) are easy and we omit their proofs.

To complete the proof, we need to show that $\alpha_{\infty} = 0$. We argue by contradiction, and suppose that $\alpha_{\infty} > 0$. Let us consider once more the function $\tilde{\varphi}_R$. Since $DJ(u_n)(\tilde{\varphi}_R u_n) \to 0$ as $n \to +\infty$, we find that

$$\beta_{\infty} \le \lambda \bar{b} \alpha_{\infty}.$$

Combining with (4.3) yields

(4.6)
$$\alpha_{\infty} \ge \left(\frac{\alpha_1}{\lambda \bar{b}}\right)^{\frac{p}{p-2}}.$$

Moreover,

$$J(u_n) - \frac{1}{2}DJ(u_n)u_n = \left(\frac{1}{2} - \frac{1}{p}\right)\lambda \int_{\mathbb{R}^N} b(x)|u_n(x)|^p dx + \left(\frac{1}{2} - \frac{1}{p}\right)\lambda \int_{\mathbb{R}^N} c(x)|u_n(x)|^q dx$$
$$\geq \frac{(p-2)\lambda}{2p} \int_{\mathbb{R}^N} b(x)|u_n(x)|^p \tilde{\varphi}_R(x) dx,$$

and letting first $n \to +\infty$ and then $R \to +\infty$ we have that

$$I_{\lambda} \ge \frac{p-2}{2p} \lambda \bar{b} \alpha_{\infty}.$$

Together with (4.6) we get the contradiction

$$I_{\lambda} \geq \frac{p-2}{2p} \left(\frac{\alpha_1}{\lambda \bar{b}}\right)^{\frac{p}{p-2}} \bar{b} \lambda = \frac{p-2}{2p} \alpha_1^{\frac{p}{p-2}} \left(\lambda \bar{b}\right)^{\frac{2}{2-p}}.$$

Remark 4.4. The lack of precise information about solutions to the limit equation

$$(I - \Delta)^{\alpha} u = \bar{b} |u|^{p-2} u \quad \text{in } \mathbb{R}^N,$$

and particularly the lack of a uniqueness result does not allow us to state the assumption (4.2) in terms of the ground state energy of the associated Euler functional.

5. Concave-convex nonlinearities and sign-changing potentials

In this section, following [5,20], we consider the equation

(5.1)
$$(I - \Delta)^{\alpha} u = \lambda b(x) |u|^{p-2} u + c(x) |u|^{q-2} u in \mathbb{R}^{N}$$

where $N > 2\alpha$, $\lambda > 0$ and and 1 . We assume that both <math>|b| and |c| satisfy the compactness conditions (K).

Remark 5.1. If $u_n \to 0$ weakly in $L^{\alpha,2}(\mathbb{R}^N)$, it follows from Proposition 3.5 that

$$\left| \int_{\mathbb{R}^N} b(x) |u_n(x)|^p \, dx \right| \le \int_{\mathbb{R}^N} |b(x)| |u_n(x)|^p \, dx \to 0$$

as $n \to +\infty$ along a subsequence. Therefore $\lim_{n \to +\infty} \int_{\mathbb{R}^N} b(x) |u_n(x)|^p dx = 0$. By the same token, $\lim_{n \to +\infty} \int_{\mathbb{R}^N} c(x) |u_n(x)|^q dx = 0$.

Definition 5.2. We say that a continuous function w changes sign if both sets $\{x \in \mathbb{R}^N \mid w(x) > 0\}$ and $\{x \in \mathbb{R}^N \mid w(x) < 0\}$ are nonempty.

Let us recall from [15] that (5.1) is equivalent to the Neumann system

(5.2)
$$\begin{cases} -\operatorname{div}(y^{1-2\alpha}\nabla v) + y^{1-2\alpha}v = 0 & \text{in } \mathbb{R}^{N+1}_+ \\ -y^{1-2\alpha}\frac{\partial v}{\partial y} = \kappa_{\alpha} \left\{ \lambda b(x) |u(x,0)|^{p-2} u(x,0) + c(x) |u(x,0)|^{q-2} u(x,0) \right\} & \text{on } \mathbb{R}^{N} \times \{0\} \end{cases}$$

Weak solutions to (5.2) correspond to critical points of the Euler functional E_{λ} : $H^{1}(\mathbb{R}^{N+1}_{+}, y^{1-2\alpha}) \to \mathbb{R}$ defined by

$$E_{\lambda}(v) = \frac{1}{2} \int_{\mathbb{R}^{N+1}_{+}} y^{1-2\alpha} \left(|\nabla v(x,y)|^{2} + v(x,y)^{2} \right) dx dy$$
$$- \frac{\lambda}{p} \int_{\mathbb{R}^{N}} b(x) |\operatorname{tr} v(x)|^{p} dx - \frac{1}{q} \int_{\mathbb{R}^{N}} c(x) |\operatorname{tr} v(x)|^{q} dx,$$

where $H^1(\mathbb{R}^{N+1}_+,y^{1-2\alpha})$ is the completion of $C_0^\infty(\mathbb{R}^{N+1}_+)$ with respect to the weighted norm

$$||v||_{H^1(\mathbb{R}^{N+1}_+, y^{1-2\alpha})} = \left(\int_{\mathbb{R}^{N+1}_+} y^{1-2\alpha} \left(|\nabla v(x, y)|^2 + v(x, y)^2 \right) \, dx \, dy \right)^{1/2}$$

and tr: $H^1(\mathbb{R}^{N+1}_+, y^{1-2\alpha}) \to L^{\alpha,2}(\mathbb{R}^N)$ is the continuous trace operator defined in [15, Proposition 6.2]. In particular

(5.3)
$$\kappa_{\alpha} \|\operatorname{tr} v\|_{L^{\alpha,2}(\mathbb{R}^N)}^2 \le \|v\|_{H^1(\mathbb{R}^{N+1})}^2 \quad \text{for all } v \in H^1(\mathbb{R}^{N+1}).$$

Remark 5.3. Up to a rescaling, we can (and will) assume that $\kappa_{\alpha} = 1$.

The Nehari manifold associated to (5.2) is

$$\mathcal{N}_{\lambda} = \left\{ v \in H^{1}(\mathbb{R}^{N+1}_{+}, y^{1-2\alpha}) \setminus \{0\} \mid DE_{\lambda}(v)v = 0 \right\}$$

or the set of those $v \in H^1(\mathbb{R}^{N+1}_+, y^{1-2\alpha}) \setminus \{0\}$ satisfying

$$\int_{\mathbb{R}^{N+1}_+} y^{1-2\alpha} \left(|\nabla v(x,y)|^2 + v(x,y)^2 \right) \, dx \, dy = \lambda \int_{\mathbb{R}^N} b(x) |\operatorname{tr} v(x)|^p \, dx + \int_{\mathbb{R}^N} c(x) |\operatorname{tr} v(x)|^q \, dx.$$

To each $v \in H^1(\mathbb{R}^{N+1}_+, y^{1-2\alpha})$ we attach its fiber map $\varphi_v \colon [0, +\infty) \to \mathbb{R}$ defined by $\varphi_v(t) = E_\lambda(tv)$. By a direct calculation,

$$D\varphi_{v}(t) = t \int_{\mathbb{R}^{N+1}_{+}} y^{1-2\alpha} \left(|\nabla v(x,y)|^{2} + v(x,y)^{2} \right) dx dy - \lambda t^{p-1} \int_{\mathbb{R}^{N}} b(x) |\operatorname{tr} v(x)|^{p} dx - t^{q-1} \int_{\mathbb{R}^{N}} c(x) |\operatorname{tr} v(x)|^{q} dx$$

and

$$D^{2}\varphi_{v}(t) = \int_{\mathbb{R}^{N+1}_{+}} y^{1-2\alpha} \left(|\nabla v(x,y)|^{2} + v(x,y)^{2} \right) dx dy - (p-1)\lambda t^{p-2} \int_{\mathbb{R}^{N}} b(x) |\operatorname{tr} v(x)|^{p} dx - (r-1)t^{q-2} \int_{\mathbb{R}^{N}} c(x) |\operatorname{tr} v(x)|^{q} dx$$

Remark 5.4. Clearly, $tv \in \mathcal{N}_{\lambda}$ if and only if $D\varphi_v(t) = 0$. In particular, $v \in \mathcal{N}_{\lambda}$ if and only if $D\varphi_v(1) = 0$.

We decompose

$$\mathcal{N}_{\lambda} = \mathcal{N}_{\lambda}^{-} \cup \mathcal{N}_{\lambda}^{0} \cup \mathcal{N}_{\lambda}^{+},$$

where

$$\mathcal{N}_{\lambda}^{-} = \left\{ v \in \mathcal{N}_{\lambda} \mid D^{2} \varphi_{v}(1) < 0 \right\}$$
$$\mathcal{N}_{\lambda}^{0} = \left\{ v \in \mathcal{N}_{\lambda} \mid D^{2} \varphi_{v}(1) = 0 \right\}$$
$$\mathcal{N}_{\lambda}^{+} = \left\{ v \in \mathcal{N}_{\lambda} \mid D^{2} \varphi_{v}(1) > 0 \right\}.$$

Let us define

$$C^{+} = \left\{ v \in H^{1}(\mathbb{R}^{N+1}_{+}, y^{1-2\alpha}) \mid \int_{\mathbb{R}^{N}} c(x) |\operatorname{tr} v(x)|^{q} dx > 0 \right\}$$

$$C^{-} = \left\{ v \in H^{1}(\mathbb{R}^{N+1}_{+}, y^{1-2\alpha}) \mid \int_{\mathbb{R}^{N}} c(x) |\operatorname{tr} v(x)|^{q} dx < 0 \right\}$$

$$C^{0} = \left\{ v \in H^{1}(\mathbb{R}^{N+1}_{+}, y^{1-2\alpha}) \mid \int_{\mathbb{R}^{N}} c(x) |\operatorname{tr} v(x)|^{q} dx = 0 \right\}$$

$$B^{+} = \left\{ v \in H^{1}(\mathbb{R}^{N+1}_{+}, y^{1-2\alpha}) \mid \int_{\mathbb{R}^{N}} b(x) |\operatorname{tr} v(x)|^{p} dx > 0 \right\}$$

$$B^{-} = \left\{ v \in H^{1}(\mathbb{R}^{N+1}_{+}, y^{1-2\alpha}) \mid \int_{\mathbb{R}^{N}} b(x) |\operatorname{tr} v(x)|^{p} dx < 0 \right\}$$

$$B^{0} = \left\{ v \in H^{1}(\mathbb{R}^{N+1}_{+}, y^{1-2\alpha}) \mid \int_{\mathbb{R}^{N}} b(x) |\operatorname{tr} v(x)|^{p} dx < 0 \right\}.$$

Lemma 5.5. (1) If $v \in B^- \cap C^-$, then no multiple of v belongs to \mathcal{N}_{λ} .

(2) If either $v \in B^+ \cap C^-$ or $v \in B^- \cap C^+$, then there exists one and only one t(v) > 0 such that $t(v)v \in \mathcal{N}_{\lambda}$.

Proof. Both statements can be proved by an elementary inspection of the fibering map φ_v . \square The case $v \in B^+ \cap C^+$ requires some additional care.

Lemma 5.6. There exists $\lambda_0 > 0$ such that $\lambda < \lambda_0$ implies $\varphi_v > 0$ for all $v \in H^1(\mathbb{R}^{N+1}_+, y^{1-2\alpha})$. If $\lambda < \lambda_0$ and $v \in B^+ \cap C^+$, then φ_v possesses exactly two critical points.

Proof. Let $v \in H^1(\mathbb{R}^{N+1}_+, y^{1-2\alpha})$ be such that $\int_{\mathbb{R}^N} c(x) |\operatorname{tr} v(x)|^q dx > 0$. We introduce

$$F_v(t) = \frac{t^2}{2} \int_{\mathbb{R}^{N+1}_+} y^{1-2\alpha} \left(|\nabla v(x,y)|^2 + v(x,y)^2 \right) dx dy - \frac{t^q}{q} \int_{\mathbb{R}^N} c(x) |\operatorname{tr} v(x)|^q dx.$$

Since

$$DF_v(t) = t \int_{\mathbb{R}^{N+1}_+} y^{1-2\alpha} \left(|\nabla v(x,y)|^2 + v(x,y)^2 \right) \, dx \, dy - t^{q-1} \int_{\mathbb{R}^N} c(x) |\operatorname{tr} v(x)|^q \, dx,$$

the function F_v attains its maximum at

$$t^* = \left(\frac{\int_{\mathbb{R}^{N+1}_+} y^{1-2\alpha} \left(|\nabla(x,y)|^2 + v(x,y)^2 \right) \, dx \, dy}{\int_{\mathbb{R}^N} c(x) |\operatorname{tr} v(x)|^q \, dx} \right)^{\frac{1}{q-2}}.$$

Moreover

$$F_{v}(t^{*}) = \left(\frac{1}{2} - \frac{1}{q}\right) \left(\frac{\left(\int_{\mathbb{R}^{N+1}_{+}} y^{1-2\alpha} \left(|\nabla(x,y)|^{2} + v(x,y)^{2}\right) dx dy\right)^{q}}{\left(\int_{\mathbb{R}^{N}} c(x)|\operatorname{tr} v(x)|^{q} dx\right)^{2}}\right)^{\frac{1}{q-2}}$$
$$D^{2}F_{v}(t^{*}) = (1-q) \int_{\mathbb{R}^{N+1}_{+}} y^{1-2\alpha} \left(|\nabla v(x,y)|^{2} + v(x,y)^{2}\right) dx dy < 0.$$

Let S_q be the best constant for the inequality

$$\left(\int_{\mathbb{R}^N} |\operatorname{tr} v(x)|^q \, dx \right)^{\frac{1}{q}} \le S_q \sqrt{\int_{\mathbb{R}^{N+1}_+} y^{1-2\alpha} \left(|\nabla v(x,y)|^2 + v(x,y)^2 \right) \, dx \, dy}$$

that follows from (5.3) and Theorem 2.2. Then

$$F_v(t^*) \ge \left(\frac{1}{2} - \frac{1}{q}\right) \left(\frac{1}{\|b^+\|_{\infty} S_q^{2r}}\right)^{\frac{1}{q-2}} = \delta > 0$$

with δ independent of v. Moreover,

$$\frac{(t^{*})^{p}}{p} \int_{\mathbb{R}^{N}} b(x) |\operatorname{tr} v(x)|^{p} dx
\leq \frac{\|b\|_{\infty}}{p} S_{p}^{p} \left(\frac{\int_{\mathbb{R}^{N+1}_{+}} y^{1-2\alpha} \left(|\nabla v(x,y)|^{2} + v(x,y)^{2} \right) dx dy}{\int_{\mathbb{R}^{N}} c(x) |\operatorname{tr} v(x)|^{q} dx} \right)^{\frac{p}{q-2}} \|v\|_{H^{1}(\mathbb{R}^{N+1}_{+}, y^{1-2\alpha})}^{p}
= \frac{\|b\|_{\infty}}{p} S_{p}^{p} \left(\frac{\left(\int_{\mathbb{R}^{N+1}_{+}} y^{1-2\alpha} \left(|\nabla v(x,y)|^{2} + v(x,y)^{2} \right) dx dy \right)^{q}}{\left(\int_{\mathbb{R}^{N}} c(x) |\operatorname{tr} v(x)|^{q} dx \right)^{2}} \right)^{\frac{p}{2(q-2)}}
= \frac{\|b\|_{\infty}}{p} S_{p}^{p} \left(\frac{2q}{q-2} \right)^{\frac{p}{2}} F_{v}(t^{*})^{\frac{p}{2}} = cF_{v}(t^{*})^{\frac{p}{2}}.$$

Therefore

$$\varphi_v(t^*) \ge F_v(t^*) - \lambda c F_v(t^*)^{\frac{p}{2}} = F_v(t^*)^{\frac{p}{2}} \left(F_v(t^*)^{\frac{2-p}{2}} - \lambda c \right) \ge \delta^{\frac{p}{2}} \left(\delta^{\frac{2-p}{2}} - \lambda c \right).$$

We complete the proof by choosing $\lambda < c^{-1}\delta^{\frac{2-p}{2}} = \lambda_0$.

Corollary 5.7. Let $\lambda < \lambda_0$. There exists $\delta_1 > 0$ such that $E_{\lambda}(v) \geq \delta_1$ for every $v \in \mathcal{N}_{\lambda}^-$.

Proof. Indeed, if $v \in \mathcal{N}_{\lambda}^-$, then φ_v has a positive global maximum at t=1, and in addition

$$\int_{\mathbb{R}^N} b(x) |\operatorname{tr} v(x)|^p \, dx > 0.$$

Hence

$$E_{\lambda}(v) = \varphi_v(1) = \varphi_v(t^*) \ge F_v(t^*)^{\frac{p-2}{2}} \left(F_v(t^*)^{\frac{2-p}{2}} - \lambda c \right) \ge \delta^{\frac{p}{2}} \left(\delta^{\frac{2-p}{2}} - \lambda c \right) > 0$$
 provided that $\lambda < \lambda_0$.

Corollary 5.8. If $\lambda \in (0, \lambda_0)$, then $\mathcal{N}_{\lambda}^0 = \emptyset$.

Lemma 5.9. Let v be a local minimum point of E_{λ} on either $\mathcal{N}_{\lambda}^{+}$ or $\mathcal{N}_{\lambda}^{-}$ such that $v \notin \mathcal{N}_{\lambda}^{0}$. Then $DE_{\lambda}(v) = 0$.

Proof. By the Lagrange Multiplier Rule, $DE_{\lambda}(v) = \mu DI_{\lambda}(v)$ for some $\mu \in \mathbb{R}$, where we have set $I_{\lambda}(v) = DE_{\lambda}(v)v$. Hence $DE_{\lambda}(v)v = \mu DI_{\lambda}(v)v = \mu D^{2}\varphi_{v}(1) = 0$. Since $v \notin \mathcal{N}_{\lambda}^{0}$ we must have $\mu = 0$.

Lemma 5.10. The restriction of E_{λ} to \mathcal{N} is bounded from below and coercive.

Proof. Indeed, for any $v \in \mathcal{N}_{\lambda}$, we have

$$E_{\lambda}(v) = \left(\frac{1}{2} - \frac{1}{q}\right) \|v\|_{H^{1}(\mathbb{R}^{N+1}_{+}, y^{1-2\alpha})}^{2} = \lambda \left(\frac{1}{p} - \frac{1}{q}\right) \int_{\mathbb{R}^{N}} b(x) |\operatorname{tr} v(x)|^{p} dx$$
$$\geq c_{1} \|v\|_{H^{1}(\mathbb{R}^{N+1}_{+}, y^{1-2\alpha})}^{2} - c_{2} \|v\|_{H^{1}(\mathbb{R}^{N+1}_{+}, y^{1-2\alpha})}^{p},$$

and we conclude since p < 2.

Lemma 5.11. If $\lambda < \lambda_0$, then E_{λ} attains its minimum on \mathcal{N}_{λ}^+ .

Proof. We know from Lemma 5.10 that E_{λ} is bounded from below on $\mathcal{N}_{\lambda}^{+}$. Therefore there exists a sequence $\{v_{k}\}_{k} \subset \mathcal{N}_{\lambda}^{+}$ with

$$\lim_{k \to +\infty} E_{\lambda}(v_k) = \inf_{\mathcal{N}_{\lambda}^+} E_{\lambda} > -\infty.$$

Again from Lemma 5.10 the sequence $\{v_k\}_k$ is bounded. We are then allowed to suppose, without loss of generality, that it converges weakly to some v_λ in $H^1(\mathbb{R}^{N+1}_+, y^{1-2\alpha})$. Let us pick $v \in H^1(\mathbb{R}^{N+1}_+, y^{1-2\alpha})$ such that $\int_{\mathbb{R}^N} b(x) |\operatorname{tr} v(x)|^p dx > 0$; then by Lemma 5.5 there exists $t_1 = t_1(v) > 0$ such that $t_1v \in \mathcal{N}^+_\lambda$ and $E_\lambda(t_1v) < 0$. This implies that $\inf_{\mathcal{N}^+_\lambda} E_\lambda < 0$. On \mathcal{N}_λ we have

$$E_{\lambda}(v_k) = \left(\frac{1}{2} - \frac{1}{q}\right) \|v_k\|_{H^1(\mathbb{R}^{N+1}_+, y^{1-2\alpha})}^2 - \lambda \left(\frac{1}{p} - \frac{1}{q}\right) \int_{\mathbb{R}^N} b(x) |\operatorname{tr} v_k(x)|^p dx$$

and thus

$$\lambda \left(\frac{1}{p} - \frac{1}{q} \right) \int_{\mathbb{R}^N} b(x) |\operatorname{tr} v_k(x)|^p dx = \left(\frac{1}{2} - \frac{1}{q} \right) \|v_k\|_{H^1(\mathbb{R}^{N+1}_+, y^{1-2\alpha})}^2 - E_{\lambda}(v_k).$$

By Proposition 3.5, Remark 5.1 and a semicontinuity argument, we let $k \to +\infty$ and get

$$\lambda \left(\frac{1}{p} - \frac{1}{q} \right) \int_{\mathbb{R}^N} b(x) |\operatorname{tr} v(x)|^p \, dx \ge \left(\frac{1}{2} - \frac{1}{q} \right) ||v||_{H^1(\mathbb{R}^{N+1}_+, y^{1-2\alpha})}^2 - \inf_{\mathcal{N}^+_{\lambda}} E_{\lambda} > 0.$$

We finally claim that $v_k \to v_\lambda$ strongly. If not, then

$$\int_{\mathbb{R}^{N+1}_+} y^{1-2\alpha} \left(|\nabla v(x,y)|^2 + v(x,y)^2 \right) \, dx \, dy < \liminf_{k \to +\infty} \int_{\mathbb{R}^{N+1}_+} y^{1-2\alpha} \left(|\nabla v_k(x,y)|^2 + v_k(x,y)^2 \right) \, dx \, dy.$$

Let us recall that

$$D\varphi_{v_k}(1) = t \int_{\mathbb{R}^{N+1}_+} y^{1-2\alpha} \left(|\nabla v_k(x,y)|^2 + v_k(x,y)^2 \right) dx dy$$
$$-\lambda t^{p-1} \int_{\mathbb{R}^N} b(x) |\operatorname{tr} v_k(x)|^p dx - t^{q-1} \int_{\mathbb{R}^N} c(x) |\operatorname{tr} v_k(x)|^q dx$$

and

$$D\varphi_{v_{\lambda}}(1) = t \int_{\mathbb{R}^{N+1}_{+}} y^{1-2\alpha} \left(|\nabla v_{\lambda}(x,y)|^{2} + v_{\lambda}(x,y)^{2} \right) dx dy$$
$$-\lambda t^{p-1} \int_{\mathbb{R}^{N}} b(x) |\operatorname{tr} v_{\lambda}(x)|^{p} dx - t^{q-1} \int_{\mathbb{R}^{N}} c(x) |\operatorname{tr} v_{\lambda}(x)|^{q} dx.$$

Let $t_{\lambda} = t_{\lambda}(v_{\lambda}) > 0$ be chosen in order that $t_{\lambda}v_{\lambda} \in \mathcal{N}_{\lambda}^{+}$: this is possible by Lemma 5.5. For $k \gg 1$, we have $D\varphi_{v_{k}}(t_{\lambda}) > 0$, which yields $t_{\lambda} > 1$. But then, by definition of t_{λ} ,

$$E_{\lambda}(t_{\lambda}v_{\lambda}) < E_{\lambda}(v_{\lambda}) < \lim_{k \to +\infty} E_{\lambda}(v_{k}) = \inf_{\mathcal{N}_{\lambda}^{+}} E_{\lambda}.$$

This is impossible, and we must have $v_k \to v_\lambda$ strongly in $H^1(\mathbb{R}^{N+1}_+, y^{1-2\alpha})$ as $k \to +\infty$. Since $\mathcal{N}^0_\lambda = \emptyset$, there results $v_\lambda \in \mathcal{N}^+$, and the proof is complete.

Lemma 5.12. If $\lambda < \lambda_0$, then E_{λ} attains its minimum on \mathcal{N}_{λ}^- .

Proof. We know from Corollary 5.7 that $\inf_{\mathcal{N}_{\lambda}^{-}} E_{\lambda} > 0$. Let $\{v_k\}_k \subset \mathcal{N}_{\lambda}^{-}$ be a minimizin sequence, i.e.

$$\lim_{k \to +\infty} E_{\lambda}(v_k) = \inf_{\mathcal{N}_{\lambda}^-} E_{\lambda} > 0.$$

By coercivity on $\mathcal{N}_{\lambda}^{-}$ we may assume without loss of generality that $\{v_k\}_k$ converges weakly to some v_{λ} in $H^1(\mathbb{R}^{N+1}_+, y^{1-2\alpha})$. Now,

$$E_{\lambda}(v_k) = \left(\frac{1}{2} - \frac{1}{p}\right) \|v_k\|_{H^1(\mathbb{R}^{N+1}_+, y^{1-2\alpha})}^2 + \left(\frac{1}{p} - \frac{1}{q}\right) \int_{\mathbb{R}^N} c(x) |\operatorname{tr} v_k(x)|^q dx.$$

By the assumptions on c, Proposition 3.5 and Remark 5.1

$$\lim_{k \to +\infty} \int_{\mathbb{R}^N} c(x) |\operatorname{tr} v_k(x)|^q dx = \int_{\mathbb{R}^N} c(x) |\operatorname{tr} v_\lambda(x)|^q dx.$$

Hence $\int_{\mathbb{R}^N} c(x) |\operatorname{tr} v_{\lambda}(x)|^q dx > 0$. As a consequence by Lemma 5.5, the fiber function $\varphi_{v_{\lambda}}$ possesses a global maximum at some $\tilde{t} = \tilde{t}(v_l)$ such that $\tilde{t}v_{\lambda} \in \mathcal{N}_{\lambda}^-$. On the other hand, $v_k \in \mathcal{N}_{\lambda}^-$ already implies that 1 is a global maximum point for φ_{v_k} , namely $\varphi_{v_k}(t) \leq \varphi_{v_k}(1)$ for every t > 0.

If $\{v_k\}_k$ does not converge strongly to v_λ in $H^1(\mathbb{R}^{N+1}_+, y^{1-2\alpha})$, then

$$E_{\lambda}(\tilde{t}v_{\lambda}) = \frac{\tilde{t}^{2}}{2} \int_{\mathbb{R}^{N+1}_{+}} y^{1-2\alpha} \left(|\nabla v_{\lambda}|^{2} + v_{\lambda}^{2} \right) dx dy - \frac{\lambda \tilde{t}^{p}}{p} \int_{\mathbb{R}^{N}} b(x) |\operatorname{tr} v_{\lambda}(x)|^{p} dx$$

$$- \frac{\tilde{t}^{q}}{q} \int_{\mathbb{R}^{N}} c(x) |\operatorname{tr} v_{\lambda}(x)|^{q} dx$$

$$< \liminf_{k \to +\infty} \frac{\tilde{t}^{2}}{2} \int_{\mathbb{R}^{N+1}_{+}} y^{1-2\alpha} \left(|\nabla v_{k}|^{2} + v_{k}^{2} \right) dx dy - \frac{\lambda \tilde{t}^{p}}{p} \int_{\mathbb{R}^{N}} b(x) |\operatorname{tr} v_{k}(x)|^{p} dx$$

$$- \frac{\tilde{t}^{q}}{q} \int_{\mathbb{R}^{N}} c(x) |\operatorname{tr} v_{k}(x)|^{q} dx$$

$$\leq \lim_{k \to +\infty} E_{\lambda}(\tilde{t}v_{k}) \leq \lim_{k \to +\infty} E_{\lambda}(v_{k}) = \inf_{\mathcal{N}^{-}} E_{\lambda},$$

a contradiction. We have proved that $v_k \to v_\lambda$ strongly in $H^1(\mathbb{R}^{N+1}_+, y^{1-2\alpha})$ as $k \to +\infty$, and $\mathcal{N}^0_\lambda = \emptyset$ finally implies that $v_\lambda \in \mathcal{N}^-_\lambda$.

We are ready to prove our main result.

Theorem 5.13. Let $N > 2\alpha$ and 1 . Suppose that both <math>|b| and |c| satisfy the compactness conditions (K) and that b and c change sign according to Definition 5.2. Then there exists $\lambda_0 > 0$ such that for every $0 < \lambda < \lambda_0$ equation (5.1) has at least two non-negative solutions.

Proof. We only have to prove that the two minima of E_{λ} on $\mathcal{N}_{\lambda}^{+}$ and on $\mathcal{N}_{\lambda}^{-}$ respectively can be chosen to be non-negative. To this aim, we let

$$f(x,t) = b(x)|t|^{p-2}t + c(x)|t|^{q-2}t$$
$$f_{+}(x,t) = \begin{cases} f(x,t) & \text{if } t \ge 0\\ 0 & \text{if } t < 0 \end{cases}$$
$$F_{+}(x,t) = \int_{0}^{t} f_{+}(x,s) \, ds.$$

To the functional

$$E_{\lambda}^{+} : v \mapsto \|v\|_{H^{1}(\mathbb{R}^{N+1}_{+}, y^{1-2\alpha})}^{2} - \int_{\mathbb{R}^{N}} F_{+}(x, v(x)) dx$$

we can apply all the previous lemma, and thus for every $0 < \lambda < \lambda_0$ the functional E_{λ}^+ possesses two critical points $v_{\lambda,1} \in \mathcal{N}_{\lambda}^+$ and $v_{\lambda,2} \in \mathcal{N}_{\lambda}^-$. Now, denoting by $v_{\lambda,1}^-$ the negative part of $v_{\lambda,1}$, we have

$$0 = DE_{\lambda}^{+}(v_{\lambda,1})v_{\lambda,1}^{-} = \int_{\mathbb{R}^{N+1}_{+}} y^{1-2\alpha} \left(\nabla v_{\lambda,1} \cdot \nabla v_{\lambda,1}^{-} + v_{\lambda,1}v_{\lambda,1}^{-} \right) dx dy$$

$$\geq \int_{\mathbb{R}^{N+1}_{+}} y^{1-2\alpha} \left(|\nabla v_{\lambda,1}^{-}|^{2} + |v_{\lambda,1}^{-}|^{2} \right) dx dy.$$

Therefore $v_{\lambda,1} = v_{\lambda,1}^+ \ge 0$. Since the same computation holds with $v_{\lambda,2}$ in place of $v_{\lambda,1}$, the proof is complete.

6. Perspectives and open problems

As we have seen, the Bessel operator $(I-\Delta)^{\alpha}$ shares many features with the standard fractional Laplacian $(-\Delta)^{\alpha}$. However, the most striking difference between these two operators is that the latter has some scaling properties that the former does not have. As should be clear from (2.1), (2.2) and (2.3), the Bessel operator is not compatible with the semigroup \mathbb{R}_+ acting on functions as $s \star u \colon x \mapsto u(s^{-1}x)$ for s > 0. In simpler words, the Bessel operator does not scale.

In our opinion, a very challenging problem is that of finding solutions to the fractional scalar field equation

$$(6.1) (I - \Delta)^{\alpha} u = g(u) in \mathbb{R}^N$$

under the so-called *Berestycki–Lions assumptions* on g. In [3, 4] the case $\alpha = 1$ was studied under very mild assumptions on the nonlinear term g. In particular, no Ambrosetti–Rabinowitz assumption must be imposed, and no monotonicity assumption like in (f3).

However, the main tool used in [3,4] consists in a clever exploitation of the semigroup action $s \star u \colon x \mapsto u(s^{-1}x)$ for s > 0. Indeed, solutions are constructed (roughly speaking) by solving the minimization problem

$$\inf \left\{ \int_{\mathbb{R}^N} |\nabla u|^2 + u^2 \mid \int_{\mathbb{R}^N} G(u) = 1 \right\},$$

where G is the antiderivative of g. Then a rescaling from u to a suitable $s \star u$ produces a solution of (6.1) by absorbing the Lagrange multiplier. Let us propose a closely related question by considering the special case of equation (1.1)

$$(6.2) (I - \Delta)^{\alpha} u = |u|^{p-2} u in \mathbb{R}^{N}.$$

We continue to assume that 1 . In [17] it is shown that (6.2) has infinitely many solutions. We claim that it possesses a radially symmetric ground state. Indeed, we consider the quotient

$$S = \inf \left\{ \int_{\mathbb{R}^N} \left| \left(I - \Delta \right)^{\alpha/2} u \right|^2 dx \mid u \in L^{\alpha, 2}(\mathbb{R}^N), \int_{\mathbb{R}^N} |u(x)|^p dx = 1 \right\}$$

To overcome the lack of compactness in \mathbb{R}^N we can work in the subspace $L^{\alpha,2}_{\mathrm{rad}}(\mathbb{R}^N)$ consisting of radially symmetric elements of $L^{\alpha,2}(\mathbb{R}^N)$. The fact that S is attained is now an immediate consequence of the compact embedding of $L^{\alpha,2}_{\mathrm{rad}}(\mathbb{R}^N)$ into $L^p(\mathbb{R}^N)$, see [32]. To prove that the minimizer of S can be chosen to be radially symmetric, we use the following inequality: it is probably known, but we could not find a precise recerence in the literature.

Proposition 6.1. Assume that $u \in L^{\alpha,2}(\mathbb{R}^N)$, and let u^* the usual symmetric-decreasing rearrangement of u. Then

(6.3)
$$\int_{\mathbb{R}^N} \left| (I - \Delta)^{\alpha/2} u^* \right|^2 dx \le \int_{\mathbb{R}^N} \left| (I - \Delta)^{\alpha/2} u \right|^2 dx$$

Proof. We adapt the proof given by [25, Appendix A], see also [24]. By [15, Eq. (7.5)],

$$\kappa_{\alpha} \int_{\mathbb{R}^N} \left| (-\Delta + 1)^{\alpha/2} u(x) \right|^2 dx = -2\alpha \int_{\mathbb{R}^N} \frac{\vartheta(\sqrt{|\xi|^2 + 1}t) - 1}{t^{2\alpha}} |\mathcal{F}u(\xi)|^2 d\xi,$$

where κ_{α} is an explicit positive constant and ϑ solves the ordinary differential equation

$$\vartheta'' + \frac{1 - 2\alpha}{t}\vartheta' - \vartheta = 0, \quad \vartheta(0) = 1.$$

Moreover,

(6.4)
$$\int_{\mathbb{R}^N} \frac{\vartheta(\sqrt{|\xi|^2 + 1}t) - 1}{t^{2\alpha}} |\mathcal{F}u(\xi)|^2 d\xi = t^{-2\alpha} \int_{\mathbb{R}^N} \left(u(x)P(t, \cdot) * u(x) - u(x)^2 \right) dx,$$

where

$$P(t,x) = C_{N,\alpha}' t^{2\alpha} \left(|x|^2 + t^2 \right)^{-\frac{N+2\alpha}{4}} K_{\frac{N+2\alpha}{2}} \left(\sqrt{|x|^2 + t^2} \right),$$

 $P(t,\cdot)*u(x) = \int_{\mathbb{R}^N} P(t,x-y)u(y)\,dy$ and $C'_{N,\alpha}$ is another explicit positive constant. Since $K_{\frac{N+2\alpha}{2}}$ is positive and non-increasing (because $K_{\nu}(s) > 0$ and $K'_{\nu}(s) = -\frac{\nu}{s}K_{\nu}(s) - K_{\nu-1}(s) < 0$ for all s > 0), we can apply an inequality by Riesz [24, Eq. (3.9)] and conclude that the right-hand side of (6.4) decreases if we replace u by u^* . We achieve (6.3) by letting $t \to 0$.

The proof of the following result is now straightforward.

Theorem 6.2. Let 1 . Then equation (6.2) possesses a radially symmetric ground state.

Remark 6.3. By the results of [27], if $u = G_{2\alpha} * u^p$ with $u \in L^r(\mathbb{R}^N)$ and $r > \max\{p, N(p-1)/(2\alpha)\}$ then u is radially symmetric and decreasing about some point. The previous result is weaker but much easier to prove.

What happens if we replace the power $|u|^{p-2}u$ in (6.2) with a more general nonlinear term f(u) with subcritical growth at infinity? It is not hard to check that we can solve the minimum problem

$$\widetilde{S} = \inf \left\{ \int_{\mathbb{R}^N} \left| \left(I - \Delta \right)^{\alpha/2} u \right) \right|^2 dx \mid u \in L^{\alpha,2}(\mathbb{R}^N), \ \int_{\mathbb{R}^N} F(u(x)) dx = 1 \right\}$$

in the framework of radially symmetric functions, but getting rid of the associated Lagrange multiplier is a complicated task. This difficulty is again caused by the lack of scaling invariance for the Bessel fractional operator.

For very similar reasons, it seems that the Bessel operator $(I - \Delta)^{\alpha}$ does not produce strong variational identities like the local Laplacian or the fractional Laplacian. Since identities like Pohozaev's are a consequence of the action of the very same semigroup $s \star u$, we can imagine that some troubles arise. Indeed, Felmer *et al.* proved in [17, Proposition 5.1] that the following pointwise identity holds for every $\varphi \in \mathcal{S}(\mathbb{R}^N)$:

$$(6.5) (I - \Delta)^{\alpha} (x \cdot \nabla \varphi) = x \cdot \nabla [(I - \Delta)^{\alpha} \varphi] + 2\alpha (I - \Delta)^{\alpha} \varphi - 2\alpha (I - \Delta)^{\alpha - 1} \varphi.$$

Here the crucial remark is that such a formula for an "integration by parts" involves the *inverse* Bessel operator $(I - \Delta)^{\alpha-1}$, since $\alpha - 1 < 0$. In other words, a direct attempt to extend the Pohozaev identity leads to terms that are different in nature from those appearing in the problem.

Indeed, if $(I - \Delta)^{\alpha} u = f(u)$ and u decays sufficiently fast at infinity, we can justify the following formal computation, see [17] for a similar result:

$$\int_{\mathbb{R}^N} \langle x \mid \nabla u \rangle (I - \Delta)^{\alpha} u \, dx = \int_{\mathbb{R}^N} \langle x \mid \nabla u \rangle f(u(x)) \, dx.$$

Setting

$$(I) = \int_{\mathbb{R}^N} \langle x \mid \nabla u \rangle (I - \Delta)^{\alpha} u \, dx, \quad (II) = \int_{\mathbb{R}^N} \langle x \mid \nabla u \rangle f(u(x)) \, dx,$$

we deduce that

$$(I) = (2\alpha - N) \int_{\mathbb{R}^{N}} u(x) f(u(x)) \, dx + N \int_{\mathbb{R}^{N}} F(u(x)) \, dx - 2\alpha \int_{\mathbb{R}^{N}} u(x) (I - \Delta)^{\alpha - 1} u(x) \, dx$$

and

$$(II) = -N \int_{\mathbb{R}^N} F(u(x)) \, dx.$$

Therefore the solution u satisfies the variational identity

$$(6.6) \qquad 2\alpha \int_{\mathbb{R}^N} u(x) (I-\Delta)^{\alpha-1} u(x) \, dx = 2N \int_{\mathbb{R}^N} F(u(x)) \, dx - (N-2\alpha) \int_{\mathbb{R}^N} u(x) f(u(x)) \, dx.$$

The left-hand side of (6.6) is finite if and only if $u \in L^{\alpha-1,2}(\mathbb{R}^N)$, and this is certainly true since $0 < \alpha < 1$ and therefore $L^{\alpha,2}(\mathbb{R}^N) \subset L^{\alpha-1,2}(\mathbb{R}^N)$, see [26, Equation (7.3)]. Anyway, the identity (6.6) is structurally different than the usual Pohozaev identity in the local case $\alpha = 1$.

We observe that for the standard fractional Laplacian $(-\Delta)^{\alpha}$ the analogous of (6.5) reads

$$(-\Delta)^{\alpha} \langle x \mid \nabla u \rangle = 2\alpha (-\Delta)^{\alpha} u + \langle x \mid \nabla (-\Delta)^{\alpha} u \rangle,$$

see [31, page 602], so that we get a more useful variational identity.

It would be interesting to understand if classical results like those proved in [3,4] using variational identities can be somehow extended to Bessel fractional operators.

REFERENCES

- D. R. Adams and L. I. Hedberg, Function spaces and potential theory, Grundlehren der Mathematischen Wissenschaften [Fundamental Principles of Mathematical Sciences], vol. 314, Springer-Verlag, Berlin, 1996. MR1411441 (97):46024)
- [2] C. O. Alves and M. A. S. Souto, Existence of solutions for a class of nonlinear Schrödinger equations with potential vanishing at infinity, J. Differential Equations 254 (2013), no. 4, 1977–1991, DOI 10.1016/j.jde.2012.11.013. MR3003299
- [3] H. Berestycki and P.-L. Lions, Nonlinear scalar field equations. I. Existence of a ground state, Arch. Rational Mech. Anal. 82 (1983), no. 4, 313–345, DOI 10.1007/BF00250555. MR695535 (84h:35054a)
- [4] ______, Nonlinear scalar field equations. II. Existence of infinitely many solutions, Arch. Rational Mech. Anal.
 82 (1983), no. 4, 347–375, DOI 10.1007/BF00250556. MR695536 (84h:35054b)
- [5] K. J. Brown and T.-F. Wu, A fibering map approach to a semilinear elliptic boundary value problem, Electron.
 J. Differential Equations (2007), No. 69, 9. MR2308869 (2008a:35101)
- [6] C. Bucur and E. Valdinoci, Nonlocal diffusion and applications, available at arXiv:1504.08292.
- [7] R. Carmona, W. C. Masters, and B. Simon, Relativistic Schrödinger operators: asymptotic behavior of the eigenfunctions, J. Funct. Anal. 91 (1990), no. 1, 117–142, DOI 10.1016/0022-1236(90)90049-Q. MR1054115 (91i:35139)
- [8] G. Cerami, Un criterio di esistenza per i punti critici su varietà illimitate, Rend. Acad. Sci. Lett. Istituto Lombardo 112 (1978), 332–336.
- [9] J. Chabrowski, Concentration-compactness principle at infinity and semilinear elliptic equations involving critical and subcritical Sobolev exponents, Calc. Var. Partial Differential Equations 3 (1995), no. 4, 493–512, DOI 10.1007/BF01187898. MR1385297 (97j:35029)
- [10] X. Chang and Z.-Q. Wang, Ground state of scalar field equations involving a fractional Laplacian with general nonlinearity, Nonlinearity 26 (2013), no. 2, 479–494, DOI 10.1088/0951-7715/26/2/479. MR3007900

- [11] S. Cingolani and S. Secchi, Ground states for the pseudo-relativistic Hartree equation with external potential, Proc. Roy. Soc. Edinburgh Sect. A 145 (2015), no. 1, 73–90, DOI 10.1017/S0308210513000450. MR3304576
- [12] _____, Semiclassical analysis for pseudorelativistic Hartree equations, J. Differential Equations 258 (2015), 4156–4179, DOI 10.1016/j.jde.2015.01.029.
- [13] V. Coti Zelati and M. Nolasco, Existence of ground states for nonlinear, pseudo-relativistic Schrödinger equations, Atti Accad. Naz. Lincei Cl. Sci. Fis. Mat. Natur. Rend. Lincei (9) Mat. Appl. 22 (2011), no. 1, 51–72, DOI 10.4171/RLM/587. MR2799908 (2012d:35346)
- [14] ______, Ground states for pseudo-relativistic Hartree equations of critical type, Rev. Mat. Iberoam. 29 (2013), no. 4, 1421–1436, DOI 10.4171/RMI/763. MR3148610
- [15] M. M. Fall and V. Felli, Unique continuation properties for relativistic Schrödinger operators with a singular potential, available at arXiv:1312.6516.
- [16] P. Felmer, A. Quaas, and J. Tan, Positive solutions of the nonlinear Schrödinger equation with the fractional Laplacian, Proc. Roy. Soc. Edinburgh Sect. A 142 (2012), no. 6, 1237–1262, DOI 10.1017/S0308210511000746. MR3002595
- [17] P. Felmer and I. Vergara, Scalar field equations with non-local diffusion, NoDEA (2015), to appear.
- [18] R. L Frank, E. Lenzmann, and L. Silvestre, Uniqueness of radial solutions for the fractional Laplacian (2013), available at arXiv:1302.2652.
- [19] J. E. Galé, P. J. Miana, and P. R. Stinga, Extension problem and fractional operators: semigroups and wave equations, J. Evol. Equ. 13 (2013), no. 2, 343–368, DOI 10.1007/s00028-013-0182-6. MR3056307
- [20] S. Goyal and K. Sreenadh, A Nehari manifold for non-local elliptic operator with concave-convex non-linearities and sign-changing weight function, available at arXiv:1307.5149.
- [21] L. Jeanjean and K. Tanaka, A remark on least energy solutions in \mathbb{R}^N , Proc. Amer. Math. Soc. **131** (2003), no. 8, 2399–2408 (electronic), DOI 10.1090/S0002-9939-02-06821-1. MR1974637 (2004c:35127)
- [22] N. Laskin, Fractional quantum mechanics and Lévy path integrals, Phys. Lett. A 268 (2000), no. 4-6, 298–305, DOI 10.1016/S0375-9601(00)00201-2. MR1755089 (2000m:81097)
- [23] N. Laskin, Fractional Schrödinger equation, Phys. Rev. E (3) 66 (2002), no. 5, 056108, 7, DOI 10.1103/Phys-RevE.66.056108. MR1948569 (2003k:81043)
- [24] E. H. Lieb, Existence and uniqueness of the minimizing solution of Choquard's nonlinear equation, Studies in Appl. Math. 57 (1976/77), no. 2, 93–105. MR0471785 (57 #11508)
- [25] E. H. Lieb and H.-T. Yau, The Chandrasekhar theory of stellar collapse as the limit of quantum mechanics, Comm. Math. Phys. 112 (1987), no. 1, 147–174. MR904142 (89b:82014)
- [26] P.-L. Lions and E. Magenes, Problèmes aux limites non-homogènes et applications, Vol. 1, Dunod, Paris, 1969.
- [27] L. Ma and D. Chen, Radial symmetry and monotonicity for an integral equation, J. Math. Anal. Appl. 342 (2008), no. 2, 943–949, DOI 10.1016/j.jmaa.2007.12.064. MR2445251 (2009m:35151)
- [28] ______, Radial symmetry and uniqueness for positive solutions of a Schrödinger type system, Math. Comput. Modelling 49 (2009), no. 1-2, 379–385, DOI 10.1016/j.mcm.2008.06.010. MR2480059 (2010a:35072)
- [29] G. Palatucci and A. Pisante, Improved Sobolev embeddings, profile decomposition, and concentration-compactness for fractional Sobolev spaces, Calc. Var. Partial Differential Equations 50 (2014), no. 3-4, 799–829, DOI 10.1007/s00526-013-0656-y. MR3216834
- [30] P. H. Rabinowitz, On a class of nonlinear Schrödinger equations, Z. Angew. Math. Phys. 43 (1992), no. 2, 270–291, DOI 10.1007/BF00946631. MR1162728 (93h:35194)
- [31] X. Ros-Oton and J. Serra, The Pohozaev identity for the fractional Laplacian, Arch. Ration. Mech. Anal. 213 (2014), no. 2, 587–628, DOI 10.1007/s00205-014-0740-2. MR3211861
- [32] W. Sickel and L. Skrzypczak, Radial subspaces of Besov and Lizorkin-Triebel classes: extended Strauss lemma and compactness of embeddings, J. Fourier Anal. Appl. 6 (2000), no. 6, 639–662, DOI 10.1007/BF02510700. MR1790248 (2002h:46056)
- [33] E. M. Stein, Singular integrals and differentiability properties of functions, Princeton Mathematical Series, No. 30, Princeton University Press, Princeton, N.J. 1970. MR0290095 (44 #7280)
- [34] P. R. Stinga and J. L. Torrea, Extension problem and Harnack's inequality for some fractional operators, Comm. Partial Differential Equations 35 (2010), no. 11, 2092–2122, DOI 10.1080/03605301003735680. MR2754080 (2012c:35456)
- [35] R. S. Strichartz, Analysis of the Laplacian on the complete Riemannian manifold, J. Funct. Anal. 52 (1983), no. 1, 48–79, DOI 10.1016/0022-1236(83)90090-3. MR705991 (84m:58138)
- [36] C. A. Stuart, Bifurcation for Neumann problems without eigenvalues, J. Differential Equations 36 (1980), no. 3, 391–407, DOI 10.1016/0022-0396(80)90057-1. MR576158 (81m:47089)

[37] J. Tan, Y. Wang, and J. Yang, Nonlinear fractional field equations, Nonlinear Anal. **75** (2012), no. 4, 2098–2110, DOI 10.1016/j.na.2011.10.010. MR2870902 (2012k:35585)

DIPARTIMENTO DI MATEMATICA E APPLICAZIONI, UNIVERSITÀ DEGLI STUDI DI MILANO BICOCCA, VIA COZZI 55, 20255 MILANO, ITALY

 $E ext{-}mail\ address:$ Simone.Secchi@unimib.it